$$
\frac{\mathrm{d} p \mathrm{CO}_{2}}{\mathrm{~d} t}=k_{1} \frac{\left(\mathrm{Mn}\left(\mathrm{C}_{2} \mathrm{O}_{4}\right)_{2}^{-}\right)}{\left(\mathrm{C}_{2} \mathrm{O}_{4}^{4}\right)}
$$

when oxalate ion is in excess, and by the equation

$$
\frac{\mathrm{d} p_{\mathrm{CO}}^{2}}{\mathrm{~d} t}=k_{2}\left(\mathrm{Mn}^{+++}\right)\left(\mathrm{C}_{2} \mathrm{O}_{4}^{-}\right)
$$

when oxalate ion is in deficiency and fluoride ion is present to form a complex with manganic ion. This led to the adoption of the following mechanism.

The formula of the manganic oxalate complex ion was found to be Mn$\left(\mathrm{C}_{2} \mathrm{O}_{4}\right)_{2}$ -

The influence of the ionic strength on the rate was found to be in accord with that predicted using the Brönsted hypothesis.

An explanation was proposed for the formation of peroxides when the reaction takes place in the presence of oxygen.

Pasadena, California
[Contribution from the Chemical Laboratory of the University of California]
THE HEAT CAPACITY AND ENTROPY OF CARBON MONOXIDE.
HEAT OF VAPORIZATION. VAPOR PRESSURES OF SOLID AND LIQUID. FREE ENERGY TO $5000^{\circ} \mathrm{K}$. FROM SPECTROSCOPIC DATA

By J. O. Clayton ${ }^{1}$ and W. F. Giaugue<br>Received December 28, $1931 \quad$ Publishbd July 6, 1932

Carbon monoxide gas is one of the simpler diatomic molecules. Spectroscopically the normal level is a ${ }^{1 \Sigma}$ state and it has no other electronic configurations that are appreciably occupied until temperatures far in excess of the $5000^{\circ} \mathrm{K}$. limit of our calculations are reached. While nuclear spin should not produce any appreciable effect on the observed behavior of such a molecule, it is of interest to note that the predominant isotopes of both carbon and oxygen are without nuclear spin.

In view of the molecular simplicity it might be expected that the entropy obtained from the band spectrum would agree with that obtained from the ordinary application of the third law. However, this proves not to be the case and it becomes a matter of considerable practical as well as theoretical interest to determine the extent to which similar effects may exist in other molecules. In attempting to find a definite physical explanation for the discrepancy it is desirable to have an accurate quantitative measure of the amount.

It is also our purpose to consider the equilibrium C (graphite) $+1 / 2 \mathrm{O}_{2}=$ ${ }^{1}$ Shell Research Fellow, Academic year 1929-30.

CO, combining thermal data on graphite and the accurately known heat of reaction with spectroscopic data on oxygen and carbon monoxide.

Preparation of Carbon Monoxide.-The carbon monoxide was made by the well-known method of dropping formic acid into concentrated sulfuric acid. The chemically pure reagents were boiled to remove dissolved gases and were later subjected to a vacuum in the closed preparation system. The formic acid was cooled with ice during the latter procedure to prevent too rapid evaporation. The preparation system was evacuated by means of a mercury diffusion pump and found to be vacuum tight before the reagents were introduced. The first portion of the carbon monoxide to be prepared was discarded. The carbon monoxide was passed through a $50 \%$ potassium hydroxide solution and then over phosphorus pentoxide. It was then condensed in liquid air and distilled three times; the first and last portions of each distillation being discarded.

Heat Capacity Measurements and Data.-The measurements of heat capacity were carried out in gold calorimeter II described by Giauque and Wiebe. ${ }^{2}$ The method was the same as that given by the above authors and by Giauque and Johnston. ${ }^{3}$ The discussion of the various factors affecting accuracy and the estimation of the accuracy at various temperatures has been given in the latter paper and applies to the present work. It is believed that the entropy may be computed from the results with an accuracy of two-tenths of one per cent.

Three preparations of carbon monoxide were used in the measurements. These are referred to as I, II and III.

The temperature standard was the copper-constantan thermocouple "W" of Giauque, Buffington and Schulze." Since the original comparison with the hydrogen gas thermometer it has frequently been checked against the vapor pressure of oxygen. This comparison was repeated in the course of the present experiments and it was found to agree within a few hundredths of a degree at several temperatures in the liquid oxygen range. This is within the limit of the accuracy claimed, namely, $0.05^{\circ}$.

The temperature intervals and warming rates were all measured by means of the gold resistance thermometer. In correcting for the heat effect due to vaporization into the small gas volume above the condensed gas, during a temperature rise, the density of the liquid was taken from the work of Baly and Donnan ${ }^{5}$

$$
d_{(1)}=1.1604-0.0045 T
$$

As the volume of the solid is not very important in the above correction, it could be estimated with sufficient accuracy by comparison with nitrogen.
${ }^{2}$ Giauque and Wiebe, This Journal, 50, 101 (1928).
${ }^{3}$ Giauque and Johnston, ibid., 51, 2300 (1929).
${ }^{4}$ Giauque, Buffington and Schulze, ibid., 49, 2343 (1927).
s Baly and Donnan, J. Chem. Soc., 81, 907 (1902).

The amount of material was measured by means of the five-liter volumetric apparatus described by Giauque and Johnston. ${ }^{3}$ Between series II and III an accident to the measuring bulb necessitated replacement, which was done with the same precautions described for the previous apparatus. The amount of material was determined as follows:

Number of moles $=V \times P \times \frac{D}{M} \times \frac{1}{1+K T}[1+\alpha(1-P)]$
Gas density at $0^{\circ} \mathrm{C}$.
Molecular weight $D=1.2504 \mathrm{~g} . / 1,6,7$
Coefficient of thermal expansion $\left(0-100^{\circ} \mathrm{C}\right.$.)
$M=28.000$
Coefficient of deviation from Boyle's law (per atmosphere) $\alpha=-0.0005^{9}$
Gravitational acceleration (at Berkeley) $\quad=979.973 \mathrm{~cm} . / \mathrm{sec} .{ }^{10}$
The value 4.185 absolute joules equal to 1 calorie $\left(15^{\circ}\right)$ was used in converting the electrical units to calories. The relation 1.00042 absolute joules $=1$ international joule was used in calculating the energy.

The heat capacity results are given in Table I.
Table I
Heat Capacity of Carbon Monoxide
Molecular weight 28.000

| $T$. degrees <br> absolute | $\Delta T$ | $C_{P}$, calories per <br> degree per mole | Series |
| :---: | :---: | :---: | :---: |
| 14.36 | 2.640 | 1.637 | I |
| 16.94 | 2.153 | 2.458 | I |
| 19.37 | 2.328 | 3.268 | I |
| 21.93 | 2.213 | 3.976 | I |
| 24.31 | 2.172 | 4.573 | I |
| 26.64 | 2.284 | 5.114 | I |
| 29.01 | 2.345 | 5.681 | I |
| 31.56 | 2.696 | 6.272 | I |
| 39.85 | 3.001 | 8.111 | I |
| 44.21 | 3.403 | 9.055 | I |
| 44.71 | 2.293 | 9.089 | III |
| 47.90 | 3.837 | 9.888 | I |
| 48.34 | 4.427 | 9.937 | III |
| 52.34 | 3.909 | 11.01 | I |
| 55.07 | 4.775 | 11.73 | III |
| 56.82 | 4.901 | 12.71 | I |
| 59.04 | 3.025 | 13.61 | III |
| 61.55 | Transition |  |  |
| 63.47 | 2.209 | 12.02 | II |
| 64.55 | 3.570 | 12.16 | III |
| 66.02 | 2.569 | 12.30 | I |
| 68.09 | Melting Point |  |  |

[^0]TABLE I (Concluded)
$T$. iegrees
absolute
70.02
72.17
75.47
75.80
78.78
79.06
80.61
83.39
84.66
$C_{P}$, calories per
degree per mole
14.42
14.43
14.48
14.39
14.41
14.48
14.50
14.40
14.45

Series
I
II
II
I
II
I
I
I
I

Series I, 2.7594 moles. Series II, 2.8414 moles. Series III, 2.5830 moles.
The heat capacity of carbon monoxide has been measured by Eucken ${ }^{11}$ and by Clusius. ${ }^{12}$

The results of Clusius agree quite well with the present investigation, the largest deviations being below $20^{\circ} \mathrm{K}$. and near the boiling point. The average deviation is about $3 \%$. However, the algebraic deviation is only about $0.1 \%$, showing agreement as to total energy input over the region. The earlier work of Eucken does not agree as well.


Fig. 1.-Capacity of carbon monoxide in calories per degree per mole.

The results given in Table I are shown graphically in Fig. 1.
Transition and Melting Point Temperatures.-The transition and melting point temperatures were investigated with various percentages trans-
${ }^{11}$ Eucken, Ber., 18, 4 (1916).
${ }^{12}$ Clusius, Z. physik. Chem., [B] 3, 41 (1929).
formed. The transition temperature was observed over a period of nineteen hours, the melting point for eighteen hours with preparation I and for sixteen hours with preparation II. The data are collected in Table II.

Table II

| Transition temperature |  |  | Melting point temperature |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Percentage transformed | T. ${ }^{\circ} \mathrm{K}$. | Series | Percentage transformed | $T,{ }^{\circ} \mathrm{K}$. | Series |
| 15 | 61.55 | III | 3 | 68.09 | I |
| 45 | 61.55 | III | 3 | 68.07 | I |
| 45 | 61.55 | III | 20 | 68.09 | II |
| 80 | 61.55 | III | 55 | 68.09 | II |
| 80 | 61.56 | III | 85 | 68.09 | II |
| Average | 61.55 |  | Average | 68.09 |  |

In Table III the data of previous observers have been collected.
Table III
Summary of Available Transition and Melting Point Temperature Data Transition Temperature

| $T,{ }^{\circ} \mathrm{K}$. Transition Temperature |  |
| :---: | :---: |
| 60.4 | (1916) Eucken ${ }^{11}$ |
| 61.51 | (1929) Clusius ${ }^{12}$ |
| $61.55=0.05$ | This research |
|  | Melting Point |
| 62 | (1877) Cailletet ${ }^{13}$ |
| 66 | (1885) Olszewsky ${ }^{14}$ |
| 74 | (1885) Von Wroblewsky ${ }^{15}$ |
| 67.3 | (1916), Eucken ${ }^{11}$ |
| 68.22 | (1929) Clusius ${ }^{12}$ |
| 68.06 | (1931) Verschoyle ${ }^{16}$ |
| $68.09 \pm 0.05$ | This research |

Purity of Carbon Monoxide.-One of the best methods of estimating an impurity which is soluble in a liquid but insoluble in the solid phase is by calorimetric observation of the premelting effect. The short temperature range of the solid modification stable at the melting point was a slight handicap to this method in the case of carbon monoxide. However, no abnormal rise is apparent in the last point observed below the melting point. The evidence is sufficient to determine that the impurity is not greater than one-thousandth of one mole per cent. The constancy of the melting point confirms the purity although this method is greatly inferior in sensitivity to the premelting method.

Heats of Fusion and Transition.-The heats of fusion and transition were measured by starting the energy input a little below the transition or

[^1]fusion temperature, and ending a little above. A correction was made for $\int C_{P} \mathrm{~d} T$. The results are given in Table IV with those of other observers.

Table IV
Heats of Fusion and Transition of Carbon Monoxide Molecular weight 28.000
Calories per mole
152.3
150.4
$151.3 \pm 1$
144.1
151.2
$\Delta H$ fusion
199.6
199.7
199.7
$199.7=0.2$
224.1
198.2
201.5

## Series III

Series I
(1931) Average this research
(1916) Eucken ${ }^{11}$
(1929) Clusius ${ }^{12}$

Series III
Series I
Series II
(1931) Average this research
(1916) Eucken ${ }^{11}$
(1929) Eucken (corrected by Clusius) ${ }^{12}$
(1929) Clusius ${ }^{12}$

It may be noticed that the two determinations of the heat of transition differ by an amount far beyond the ordinary limit of our calorimetric error. This point will be considered later.

Vapor Pressures of Solid and Liquid Carbon Monoxide.—Vapor pressures were measured on the same material used in Series II of the calorimetric measurements but not at the same time as it was undesirable to have the manometer volume connected during calorimetric measurements. The apparatus was essentially the same as that described by Giauque, Johnston and Kelley. ${ }^{17}$ A Gaertner cathetometer with an accuracy of 0.05 mm . of mercury was used to compare the large diameter manometer with a standard meter.

The results have been represented by equations 1 and 2 . For solid carbon monoxide from the transition point at $61.55^{\circ}$ to the melting point at $68.09^{\circ}$

$$
\begin{equation*}
\log P_{(\mathrm{cm} .)}=-\frac{425.1}{T}+7.82259-0.0075960 T \tag{1}
\end{equation*}
$$

For liquid carbon monoxide

$$
\begin{equation*}
\log P_{(\mathrm{om} .)}=-\frac{477.3}{T}+11.23721-0.064129 T+2.5911 \times 10^{-4} T^{2} \tag{2}
\end{equation*}
$$

In addition an equation was calculated for the vapor pressure of the solid stable below $61.55^{\circ} \mathrm{K}$. This was done with the assistance of the various calorimetric data, including the heat of vaporization to be given below. In correcting the measured heat of vaporization to the value which would be obtained for evaporation to a dilute gas the thermodynamic equation
${ }^{17}$ Giauque, Johnston and Kelley, This Journal, 49, 2367 (1927).

$$
\left(\frac{\partial H}{\partial P}\right)_{T}=V-T\left(\frac{\partial V}{\partial T}\right)_{P}
$$

was combined with Berthelot's equation.
The heat of sublimation of the solid stable at or below $61.55^{\circ} \mathrm{K}$. was found to be 1963 calories per mole at $61.55^{\circ} \mathrm{K}$. The heat capacity of the solid can be represented very roughly but with sufficient accuracy from $20^{\circ}$ to the transition point by

$$
C_{P}=-1.25+0.24 T
$$

These several data yield the equation, for solid carbon monoxide below $61.55^{\circ} \mathrm{K}$.

$$
\begin{equation*}
\log P_{(\mathrm{om} .)}=-\frac{418.2}{T}+4.127 \log T+1.47365-0.02623 T \tag{3}
\end{equation*}
$$

Table V
Vapor Pressure of Carbon Monoxide

| $T{ }^{\circ} \mathrm{K}$. observed | $P$ int. cm. observed | $\stackrel{\Delta P}{\text { obs.-caled. }}$ | $\Delta T$ <br> obs. calcd. | Remarks |
| :---: | :---: | :---: | :---: | :---: |
| 60.24 | 1.998 | 0.019 | -0.03 | Solid |
| 61.55 | 2.819 |  |  | Transition point |
|  | 2.810 |  |  | Transition point |
|  | 2.807 |  |  | Transition point |
|  | 2.808 |  |  | Transition point |
|  | 2.810 |  |  | Transition point |
| Average | $2.811=0.005$ | 0.000 | 0.00 | Transition point |
| 62.22 | 3.275 | -. 017 | . 02 | Solid |
| 63.30 | 4.230 | . 006 | -. 01 | Solid |
| 64.31 | 5.283 | . 008 | . 01 | Solid |
| 65.28 | 6.546 | . 013 | -. 01 | Solid |
| 66.12 | 7.769 | -. 015 | . 01 | Solid |
| 67.02 | 9.325 | -. 019 | . 01 | Solid |
| 68.09 | 11.529 |  |  | Melting point |
|  | 11.531 |  |  | Melting point |
|  | 11.531 |  |  | Melting point |
| A verage | $11.530=0.005$ | 0.000 | 0.00 | Melting point |
| 69.101 | 13.651 | $-.020$ | . 009 | Liquid |
| 69.892 | 15.558 | . 008 | $-.003$ | Liquid |
| 70.690 | 17.641 | -. 007 | . 003 | Liquid |
| 72.227 | 22.364 | . 040 | -. 012 | Liquid |
| 73.246 | 25.971 | . 045 | -. 012 | Liquid |
| 73.860 | 28.325 | . 014 | -. 003 | Liquid |
| 74.586 | 31.347 | -. 003 | . 001 | Liquid |
| 75.434 | 35.229 | . 014 | -. 003 | Liquid |
| 76.351 | 39.784 | $-.025$ | . 005 | Liquid |
| 77.267 | 44.859 | . 002 | -. 000 | Liquid |
| 78.110 | 49.934 | . 008 | -. 001 | Liquid |
| 78.948 | 55.406 | -. 002 | . 000 | Liquid |
| 79.858 | 61.871 | -. 012 | . 002 | Liquid |
| 81.102 | 71.689 | . 020 | -. 002 | Liquid |
| 82.032 | 79.747 | -. 001 | . 000 | Liquid |
| 83.132 | 90.206 | -. 001 | . 000 | Liquid |

One observation on the vapor pressure of this modification was made, namely, $P=1.998 \mathrm{~cm}$. at $60.24^{\circ} \mathrm{K}$. The observed and calculated values are included in Table V. The temperatures at which observations were made on the liquid are given to 0.001 degree as they are relatively consistent to better than 0.01 degree. The principal error appears to be in the pressure measurements, the increase in $\Delta T$, observed-calculated, at the lower temperature, being due largely to this source.
The vapor pressure of carbon monoxide has been measured by Clusius and Teske ${ }^{18}$ and their data have been compared with our equations in Table VI.

Table VI
Vapor Pressure of Carbon Monoxide
Data of Clusius and Teske Compared with Equations 1 and 2

| $T{ }^{\circ}{ }^{\circ} \mathrm{K}$. observed | $P$. int. cm. observed | $\stackrel{\Delta P}{\Delta P \text { obs. }}$ | $\begin{gathered} \Delta T \\ \text { obs. caled. } \end{gathered}$ | Remarks |
| :---: | :---: | :---: | :---: | :---: |
| 60.336 | 2.041 | -0.042 | $+0.07$ | Solid |
| 60.616 | 2.083 | -. 151 | $+.25$ | Solid |
| 61.455 | 2.459 | -. 279 | $+.41$ | Solid |
| 61.985 | 3.083 | -. 033 | $+.05$ | Solid |
| 63.974 | 4.465 | -. 453 | $+.45$ | Solid |
| 64.348 | 5.519 | + . 179 | . 15 | Solid |
| 64.892 | 6.099 | +. 090 | -. 07 | Solid |
| 65.967 | 7.755 | $+.215$ | -. 13 | Solid |
| 66.201 | 7.965 | $+.051$ | -. 03 | Solid |
| 67.171 | 9.888 | $+.255$ | -. 13 | Solid |
| 67.373 | 9.963 | -. 076 | $+.04$ | Solid |
| 68.213 | 11.708 | $-.108$ | $+.05$ | Solid |
| 69.735 | 15.184 | +. 023 | -. 01 | Liquid |
| 70.801 | 17.811 | -. 146 | $+.05$ | Liquid |
| 71.457 | 19.795 | -. 077 | +. 03 | Liquid |
| 72.338 | 22.830 | $+.135$ | -. 04 | Liquid |
| 73.350 | 26.430 | $+.112$ | -. 03 | Liquid |
| 73.409 | 26.521 | -. 021 | +. 01 | Liquid |
| 73.994 | 28.745 | $-.108$ | $+.03$ | Liquid |
| 74.690 | 31.518 | -. 285 | +.07 | Liquid |
| 74.710 | 31.897 | $+.006$ | $-.00$ | Liquid |
| 74.778 | 32.144 | -. 056 | $+.01$ | Liquid |
| 75.558 | 35.650 | -. 160 | +. 03 | Liquid |
| 75.631 | 36.180 | $+.015$ | -. 00 | Liquid |
| 76.004 | 37.995 | -. 024 | $+.00$ | Liquid |
| 76.325 | 39.691 | $+.019$ | -. 00 | Liquid |
| 76.449 | 40.783 | $+.457$ | $-.08$ | Liquid |
| 76.574 | 40.878 | -. 006 | +. 02 | Liquid |
| 77.497 | 46.482 | + . 286 | -. 05 | Liquid |
| 77.499 | 46.492 | + . 283 | -. 05 | Liquid |
| 78.265 | 51.062 | $+.154$ | -. 02 | Liquid |
| 78.828 | 54.620 | $+.024$ | $-.00$ | Liquid |
| 79.696 | 60.666 | -. 022 | $+.00$ | Liquid |
| 79.967 | 62.418 | - . 278 | +. 04 | Liquid |

[^2]While the average deviation is very much greater than in the present work, the results are well represented by our equations. In fact, equation 1 represents the data of Clusius and Teske for the liquid range with a slightly smaller average deviation than was obtained with their own equation. ${ }^{19}$

From equation 1 the boiling point was found to be $81.61 \pm 0.05^{\circ} \mathrm{K}$. The boiling point data of several observers are summarized in Table VII.

Table VII
Boiling Point of Carbon Monoxide
$\quad T,{ }^{\circ} \mathrm{K}$.
80
83
81.8
81.66
81.62
81.62
$81.61 \pm 0.05$

Observer
(1885) Von Wroblewsky ${ }^{20}$
(1885) Olszewsky ${ }^{14}$
(1902) Baly and Donnan ${ }^{5}$
(1919) Von Winning (corrected to Leiden scale) ${ }^{18}$
(1929) Clusius and Teske ${ }^{18}$
(1931) Verschoyle ${ }^{18}$

This research

Heat of Vaporization of Carbon Monoxide.-The entropy of vaporization contributes roughly one-half of the total entropy for the simple gases. For this reason it is very desirable to have a direct and accurate determination of this quantity. The experimental method has been described previously. ${ }^{2,3}$

The results are given in Table VIII, which also includes the value given by Eucken and that calculated from our vapor pressure determinations. The latter method is in general inferior to the directly determined value due to uncertain correction for gas imperfection at the higher pressures and to increasing error in $\mathrm{d} P / \mathrm{d} T$ at low pressures.

[^3]

The 1434 value for the heat of vaporization obtained from the vapor pressure measurements is of course to be given no weight in comparison with the calorimetric value. This calculation, however, increases the justification for assuming that Berthelot's equation represents the behavior of carbon monoxide. This will be assumed later in connection with the effect of gas imperfection on the entropy.

The Entropy of Carbon Monoxide.-The entropy calculation is summarized in Table IX. A Debye function with $h \nu / k=79.5$ was used to extrapolate to the absolute zero.

| Table IX |  |
| :--- | :---: |
| $\quad$ Calculation of Entropy of Carbon Monoxide |  |
| $0-11.70^{\circ} \mathrm{K}$. Debye extrapolation $h \nu / k=79.5$ | 0.458 |
| $11.70-61.55$ Graphical | 9.632 |
| Transition $151.3 / 61.55$ | 2.457 |
| $61.55-68.09$ Graphical | 1.228 |
| Fusion $199.7 / 68.09$ | 2.933 |
| $68.09-81.61$ Graphical | 2.611 |
| Vaporization $1443.6 / 81.61$ | 17.689 |
| Entropy of carbon monoxide gas at boiling point | $37.01 \pm 0.1$ |
| Correction for gas imperfection assuming Berthelot gas | $\underline{0.21}$ |
| Entropy corrected to the ideal state | $37.2 \mathrm{E} . \mathrm{U}$. |

The 0.21 E . U. correction for gas imperfection was calculated by combining Berthelot's equation with the thermodynamic equation

$$
\begin{gathered}
\left(\frac{\partial S}{\partial P}\right)_{T}=-\left(\frac{\partial V}{\partial T}\right)_{P} \text { yielding } \\
\Delta S_{\text {correotion }}=R \frac{27 T_{0}^{3} P}{32 T^{3} P_{0}}
\end{gathered}
$$

The critical data were taken from the review of such data given by Pickering ${ }^{21}$

$$
T_{\circ}=134.1^{\circ} \mathrm{K} . \quad P_{\mathrm{o}}=35 \text { atmospheres }
$$

${ }^{21}$ Pickering, Sci. Pappers Bur. Standards, 21, 608 (1926).

The correction to the ideal state is necessary in order that comparison may be made with the entropy which will now be calculated from the band spectrum of carbon monoxide.
The entropy is given by the expression

$$
\begin{gather*}
S^{\circ}=\frac{3}{2} R \ln M+\frac{5}{2} R \ln T-R \ln P-\frac{5}{2} R-7.267+R \ln Q+R T \frac{\mathrm{~d} \ln Q}{\mathrm{~d} T}  \tag{4}\\
Q=\Sigma p e^{-\epsilon / k T} \text { over all quantum levels }
\end{gather*}
$$

The energy levels of carbon monoxide may be represented by the expressions

$$
\begin{aligned}
& \epsilon_{\text {rotation }}=\left[B_{\mathrm{e}}+\alpha(v+1 / 2)\right] m^{2}+\left[D_{0}+\beta(v+1 / 2)\right] m^{4} \\
& \epsilon_{\text {vibration }}=\omega_{\mathrm{e}}(v+1 / 2)+\omega_{\mathrm{e}} X_{\mathrm{e}}(v+1 / 2)^{2}
\end{aligned}
$$

Birge ${ }^{28}$ has recalculated the data of Snow and Rideal ${ }^{24}$ and gives $B_{\mathrm{e}}=$ $1.853, \alpha=0.020, D_{\mathrm{e}}=5.418 \times 10^{-6}, \beta=-6.918 \times 10^{-8}, \omega_{\mathrm{e}}=2167.4$, $\omega_{\mathrm{e}} X_{\mathrm{e}}=-12.70$. Mulholland ${ }^{25}$ has expressed $Q$ as a simple series in terms of $B$ for rigid molecules. Giauque and Overstreet ${ }^{28}$ have given an expression which is suitable for non-rigid molecules.

This avoids the considerable labor involved in summing the individual Boltzmann factors for the rotational levels.
From the above expression for $Q$ and the derivative of its logarithm used in equation 4, the entropy of carbon monoxide is calculated and given in Table X. The values for all natural constants are those given in the "International Critical Tables."

Table X
Comparison of Experimental and Spectroscopic Entropies of Carbon Monoxide Experimedtal

| $T,{ }^{\circ} \mathrm{K}$. | Spectroscopic | Actual gas | Corrected to <br> ideal state |
| :---: | :---: | :---: | :---: |
| 81.61 | 38.318 | $37.0 \neq 0.1$ | $37.2 \mathrm{E} . \mathrm{U}$. |
| 298.1 | 47.313 | $\ldots \ldots$. | 46.2 |

The "experimental" value given for $298.1^{\circ} \mathrm{K}$. was obtained by adding the calculated increase in column 2 to 37.2 .
It is clear from the discrepancy between the spectroscopic and experimental values that some random situation remains in the solid state at liquid hydrogen temperatures. This is in addition to the small amount considered by the Debye extrapolation.

The difference between the two values is roughly $R \ln 2$. We believe that
${ }^{22}$ Giauque, This Journal, 52, 4808 (1930).
${ }^{23}$ Birge, Personal communication.
${ }^{24}$ Snow and Rideal, Proc. Roy. Soc. (London), A125, 462 (1929),
${ }^{25}$ Mulholland, Proc. Cambridge Phil. Soc., 24, 280 (1928).
${ }_{26}$ Giauque and Overstreet, This Journal, 54, 1731 (1932).
the discrepancy arises as follows. If we consider a molecule with two or more otherwise equivalent positions occupied by isotopes of a given element, it is found that such a molecule enters a crystal lattice without any particular isotopic preference. However, it is true that at extremely low temperatures the almost negligible but nevertheless real energy differences between the possible arrangements will become comparable with $k T$ and thus the crystal will attain perfect order on reversible approach to the absolute zero. It is evident even when the positions are occupied by different elements that, unless the crystal lattice has a sufficiently strong preference for a certain arrangement to impose an energy difference much larger than $k T$, the several possible orientations will exist. It is suggested that this is the case in solid carbon monoxide. This seems more probable when one considers the similar atomic sizes and that it has frequently been assumed that carbon monoxide has an electron structure similar to nitrogen. A discussion of this point has been given by Lewis. ${ }^{27}$

The electron structure, symmetrical in nitrogen, would of course be somewhat deformed by the differing charges on the carbon and oxygen kernels. Nevertheless, it seems entirely reasonable to assume that the polarity thus induced is insufficient to fix the carbon and oxygen atoms in definite positions at the temperature where the solid is formed. It appears that orderly arrangement in the above respect occurs in carbon monoxide below the temperatures investigated or, what is still more probable, that rearrangement to equilibrium is not easily possible in the solid. In this connection it may be recalled that two determinations of the heat of transition given in Table IV differed by an amount considerably beyond the ordinary limit of error. While no definite conclusion is possible, the difference could be attributed to a different degree of approach to equilibrium in the solid stable below $61.55^{\circ} \mathrm{K}$. Unfortunately the experiments were concluded before the calorimetric calculations yielded this information. It would be interesting to hold carbon monoxide at various temperatures below $61.55^{\circ}$ for considerable periods of time to investigate this effect.

The amount of entropy corresponding to a complete lack of discrimination between the carbon and oxygen ends of the molecule is $R \ln 2=$ $1.38 \mathrm{E} . \mathrm{U}$. The fact that the observed discrepancy is somewhat smaller, namely, 1.1 E. U., indicates also that some approach to order has been made and that inability to obtain equilibrium is a factor.

The above type of difficulty, while possible only in certain special cases, is perhaps the most perplexing feature that has been encountered in the practical application of the third law of thermodynamics. It is perplexing because it is by no means obvious how we are to know when it exists. The crystal under investigation may appear perfect to our present powers of
${ }^{27}$ Lewis, "Valence and the Structure of Atoms and Molecules," The Chemical Catalog Co., New York, 1923, p. 126.
observation, including $x$-rays and yet the effect may be present. Certain types of organic molecules would appear to be susceptible to this effect.

The Free Energy of Carbon Monoxide.-From the expression ${ }^{22}$

$$
\frac{F^{\circ}-E_{0}^{\circ}}{T}=-\frac{3}{2} R \ln M-\frac{5}{2} R \ln T+R \ln P+7.267-R \ln Q
$$

and the band spectra data previously given, the values of $\left(F^{\circ}-E_{0}^{\circ}\right) / T$ have been calculated. The value of $R$ was taken as 1.9869 calories per degree per mole. The results are given in Table XI. The values given in heavy type have been directly calculated. The others have been interpolated by means of a difference plot making use of the Einstein harmonic oscillator function.

| TABLE XI |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\left(F^{\circ}-E_{0}^{\circ}\right) / T$ FOR CARBON MONOXIDE |  |  |  |  |
| $T,{ }^{\circ} \mathrm{K}$. | $-\left(\frac{F^{\circ}-E_{0}^{\circ}}{T}\right)$ | $T,{ }^{\circ} \mathrm{K}$. | $-\left(\frac{F^{\circ}-E_{0}^{\circ}}{T}\right)$ | $T,{ }^{\circ} \mathrm{K}$. | $-\left(\frac{F^{\circ}-E_{0}^{\circ}}{T}\right)$ |
| 250 | $\mathbf{3 9 . 1 3 8}$ | 1200 | 50.199 | 3100 | 57.773 |
| 298.1 | 40.361 | 1250 | 50.502 | 3200 | 57.832 |
| 300 | 40.405 | 1300 | 50.793 | 3300 | 58.083 |
| 350 | 41.476 | 1400 | 51.347 | 3400 | 58.328 |
| 400 | 42.404 | 1500 | 51.867 | 3500 | 58.566 |
| 450 | 43.224 | 1600 | 52.356 | 3600 | 58.798 |
| 500 | 43.959 | 1700 | 52.819 | 3700 | 59.024 |
| 550 | 44.624 | 1800 | 53.259 | 3800 | 59.245 |
| 600 | 45.233 | 1900 | 53.678 | 3900 | 59.461 |
| 650 | 45.791 | 2000 | 54.077 | 4000 | 59.671 |
| 700 | 46.317 | 2100 | 54.458 | 4100 | 59.867 |
| 750 | 46.806 | 2200 | 54.824 | 4200 | 60077 |
| 800 | 47.264 | 2300 | 55.176 | 4300 | 60.274 |
| 850 | 47.696 | 2400 | 55.513 | 4400 | 60.466 |
| 900 | 48.106 | 2500 | 55.838 | 4500 | 60.655 |
| 950 | 48.495 | 2600 | 56.152 | 4600 | 60.839 |
| 1000 | 48.867 | 2700 | 56.455 | 4700 | 61.020 |
| 1050 | 49.221 | 2800 | 56.748 | 4800 | 61.197 |
| 1100 | 49.560 | 2900 | 57.031 | 4900 | 61.372 |
| 1150 | 49.886 | 3000 | 57.307 | 5000 | 61.542 |

Before considering the reaction

$$
\mathrm{C}_{\text {Eraphite }}+1 / 2 \mathrm{O}_{2}=\mathrm{CO}
$$

it is necessary to obtain the $\left(F^{\circ}-E_{0}^{0}\right) / T$ for graphite. The values for oxygen have been tabulated by Johnston and Walker. ${ }^{28}$

The calculation for graphite was made graphically, using the heat capacity data collected in the "I. C. T." These data extend only to $2400^{\circ} \mathrm{K}$. and are not as accurate as might be desired. However, it can be seen, by analogy with other substances, that the heat capacity curve given by the available data is not likely to be much in error. The values ob-
${ }^{28}$ Johnston and Walker, have kindly allowed us to use their data in advance of publication.
tained are given in Table XII. The values for graphite have been ex.pressed to 0.001 unit although this is considerably beyond the absolute accuracy of the heat capacity data. This is in accordance with the free energy convention which thus allows accurate differences to be obtained.

| Table XII |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| ( $F^{\circ}-E_{0}^{\circ}$ )/T for Graphite |  |  |  |  |  |
| T. ${ }^{\circ} \mathrm{K}$. | $-\left(\frac{F^{\circ}-E_{0}^{0}}{T}\right)$ | $T,{ }^{\circ} \mathrm{K} .$ | - $\left(\frac{F^{\circ}-E_{0}^{\circ}}{T}\right)$ | T, ${ }^{\circ} \mathrm{K}$. | $-\left(\frac{F^{\circ}-E_{0}^{0}}{T}\right)$ |
| 250 | 0.403 | 850 | 2.325 | 1800 | 4.941 |
| 298.1 | . 545 | 900 | 2.485 | 1900 | 5.172 |
| 300 | . 551 | 9.50 | 2.643 | 2000 | 5.394 |
| 350 | . 698 | 1000 | 2.798 | 2100 | 5.608 |
| 400 | . 854 | 1050 | 2.950 | 2200 | 5.817 |
| 450 | 1.016 | 1100 | 3.099 | 2300 | 6.021 |
| 500 | 1.180 | 1150 | 3.246 | 2400 | 6.218 |
| 550 | 1.345 | 1200 | 3.391 | 2500 | 6.411 |
| 600 | 1.510 | 1250 | 3.534 | 2600 | 6.600 |
| 650 | 1.674 | 1300 | 3.674 | 2700 | 6.783 |
| 700 | 1.838 | 1400 | 3.945 | 2800 | 6.962 |
| 750 | 2.001 | 1500 | 4.206 | 2900 | 7.137 |
| 800 | 2.164 | 1600 | 4.460 | 3000 | 7.309 |
|  |  | 1700 | 4.704 |  |  |
| $\frac{\Delta F}{T}=-R \ln K=\Delta\left(\frac{F^{\circ}-E_{0}^{\circ}}{T}\right)+\Delta E_{0}^{\circ}$ |  |  |  |  |  |

$\Delta E_{0}^{\circ}$ is given by the equation
$\Delta E_{0}^{\circ}=\Delta H^{\circ}-\left[\frac{5}{2} R T+R T^{2} \frac{\mathrm{~d} \ln Q}{\mathrm{~d} T}\right]_{\mathrm{CO}}+1 / 2\left[\frac{5}{2} R T+R T^{2} \frac{\mathrm{~d} \ln Q}{\mathrm{~d} T}\right]_{\mathrm{O}_{2}}+\left[H^{\circ}-E_{0}^{\circ}\right]_{\mathrm{c}}$
The value of $\Delta H^{\circ}$ was obtained by difference of the heats of combustion of graphite and of carbon monoxide.

For the first of these the value of Roth and Naeser ${ }^{29}$ was adopted, We have learned from Professor G. S. Parks that a personal communication from Professor Roth states that the graphite weighings were not given on a vacuum basis. The application of the correction for buoyancy leads to a value of 94,240 as the heat of combustion of 12.000 g . of $\beta$ graphite at $298.1^{\circ} \mathrm{K}$.

For the second, the work of Rossini ${ }^{30}$ gives $67,623 \pm 30$ calories as the heat of combustion of carbon monoxide at $298.1^{\circ} \mathrm{K}$.

Combining

$$
\mathrm{C}_{(\beta \text { graphite })}+1 / 2 \mathrm{O}_{2}=\mathrm{CO} \Delta H_{298 \cdot 1}^{\circ}=-26,620
$$

For carbon monoxide

$$
\begin{aligned}
\frac{5}{2} R T+R T^{2} \frac{\mathrm{~d} \ln Q}{\mathrm{~d} T} & =1480.7+590.2 \\
& =2070.9 \text { calories at } 298.1^{\circ} \mathrm{K}
\end{aligned}
$$

[^4]For oxygen

$$
\begin{aligned}
\frac{1}{2}\left[\frac{5}{2} R T+R T^{2} \frac{\mathrm{~d} \ln Q}{\mathrm{~d} T}\right] & =\frac{1}{2}[1480.7+589.6] \\
& =1035.2 \text { calories at } 298.1^{\circ} \mathrm{K} .
\end{aligned}
$$

For graphite the "I. C. T." give $H^{\circ}-E_{0}^{\circ}=251.4$ calories at $298.1^{\circ} \mathrm{K}$.
Combining the several values in equation $7, \Delta E_{0}^{\circ}=-27,404$ calories per mole of carbon monoxide.

Table XIII
Free Energy of Formation of Carbon Monoxide

| T, ${ }^{\circ} \mathrm{K}$. | $\left(\frac{F^{\circ}-E_{0}^{0}}{T}\right)$ | $-\frac{\Delta F^{\circ}}{T}$ | $K=[\mathrm{CO}] /\left[\mathrm{O}_{2}\right]^{1 / 2}$ |
| :---: | :---: | :---: | :---: |
| 298.1 | 18.776 | 110.705 | $1.58 \times 10^{24}$ |
| 300 | 18.792 | 110.139 | $1.19 \times 10^{24}$ |
| 350 | 19.182 | 97.479 | $2.03 \times 10^{21}$ |
| 400 | 19.487 | 87.997 | $1.71 \times 10^{19}$ |
| 450 | 19.733 | 80.631 | $4.21 \times 10^{17}$ |
| 500 | 19.933 | 74.746 | $2.17 \times 10^{18}$ |
| 550 | 20.096 | 69.921 | $1.92 \times 10^{15}$ |
| 600 | 20.231 | 65.904 | $2.54 \times 10^{14}$ |
| 650 | 20.339 | 62.499 | $4.58 \times 10^{13}$ |
| 700 | 20.435 | 59.584 | $1.06 \times 10^{13}$ |
| 750 | 20.510 | 57.049 | $2.95 \times 10^{12}$ |
| 800 | 20.570 | 54.825 | $9.64 \times 10^{11}$ |
| 850 | 20.619 | 52.859 | $3.58 \times 10^{11}$ |
| 900 | 20.657 | 51.106 | $1.48 \times 10^{11}$ |
| 950 | 20.687 | 49.533 | $6.71 \times 10^{10}$ |
| 1000 | 20.713 | 48.117 | $3.29 \times 10^{10}$ |
| 1050 | 20.732 | 46.831 | $1.72 \times 10^{10}$ |
| 1100 | 20.747 | 45.660 | $9.55 \times 10^{9}$ |
| 1150 | 20.757 | 44.587 | $5.57 \times 10^{9}$ |
| 1200 | 20.763 | 43.600 | $3.39 \times 10^{9}$ |
| 1250 | 20.766 | 42.689 | $2.14 \times 10^{9}$ |
| 1300 | 20.766 | 41.846 | $1.40 \times 10^{9}$ |
| 1400 | 20.762 | 40.336 | $6.56 \times 10^{8}$ |
| 1500 | 20.751 | 39.020 | $3.38 \times 10^{8}$ |
| 1600 | 20.733 | 37.861 | $1.89 \times 10^{8}$ |
| 1700 | 20.712 | 36.832 | $1.13 \times 10^{8}$ |
| 1800 | 20.688 | 35.912 | $7.08 \times 10^{7}$ |
| 1900 | 20.659 | 35.082 | $4.66 \times 10^{7}$ |
| 2000 | 20.629 | 34.331 | $3.19 \times 10^{7}$ |
| 2100 | 20.599 | 33.649 | $2.26 \times 10^{7}$ |
| 2200 | 20.568 | 33.024 | $1.65 \times 10^{7}$ |
| 2300 | 20.534 | 32.449 | $1.24 \times 10^{7}$ |
| 2400 | 20.500 | 31.918 | $9.48 \times 10^{6}$ |
| 2500 | 20.465 | 31.427 | $7.40 \times 10^{8}$ |
| 2600 | 20.428 | 30.968 | $5.88 \times 10^{6}$ |
| 2700 | 20.392 | 30.542 | $4.74 \times 10^{6}$ |
| 2800 | 20.355 | 30.142 | $3.87 \times 10^{8}$ |
| 2900 | 20.317 | 29.767 | $3.21 \times 10^{8}$ |
| 3000 | 20.279 | 29.414 | $2.69 \times 10^{6}$ |

From this value, the several values of $\left(F^{\circ}-E_{0}^{\circ}\right) / T$ and equation 6 , $\Delta F^{o} / T$ and $K$ have been computed for the formation of carbon monoxide from the elements. The results are given in Table XIII.

The values of $K$ should be accurate to about $5 \%$. At the lower temperatures the error is chiefly in $\Delta E_{0}^{\circ}$ and at the higher temperatures the uncertainty is almost entirely due to the data on graphite.
In the absence of better information on the heat capacity of carbon above $3000^{\circ}$ it seems preferable to omit values above $3000^{\circ}$. However, the values for carbon monoxide have been extended to $5000^{\circ}$ and the values for $\left(F^{\circ}-E_{0}^{\circ}\right) / T$ of oxygen given by Johnston and Walker ${ }^{28}$ extend to this temperature. Thus by extrapolation of the data for graphite above $3000^{\circ} \mathrm{K}$. a quite reliable calculation of the equilibrium is available.

We wish to thank H. L. Johnston for assisting with the measurements of series I and R. W. Blue for assisting with series II and III.

## Summary

The heat capacities of the two crystalline forms of carbon monoxide and the liquid have been measured from $13^{\circ} \mathrm{K}$. to the boiling point.

The transition temperature was found to be $61.55^{\circ} \mathrm{K} . \pm 0.05$, the melting point $68.09^{\circ} \mathrm{K} . \pm 0.05$ and the boiling point $81.61^{\circ} \mathrm{K} . \pm 0.05$. From the calorimetric measurements the heat of transition was determined as 151.3 $\pm 1$ calories per mole, the heat of fusion $199.7 \pm 0.2$ and the heat of vaporization $1443.6 \pm 1.0$ calories per mole, measured at 760 mm .

From the experimental data the entropy of carbon monoxide gas was calculated to be $37.0 \pm 0.1 \mathrm{E}$. U. at $81.61^{\circ} \mathrm{K}$. Assuming Berthelot's gas equation a correction amounting to 0.2 E . U . for gas imperfection was added, giving 37.2 E . U. This value can be compared to the value calculated for the ideal gas from the band spectrum, 38.318 at $81.61^{\circ} \mathrm{K}$.
The observed entropy is low by an amount corresponding approximately to $R \ln 2$. Reasons are given for believing that this is due to lack of discrimination within the crystal lattice for the oxygen and carbon ends of the molecule. It is suggested that this effect may exist in other molecules and that a quantitative correction will be possible in many cases.
The entropy of carbon monoxide gas at $298.1^{\circ} \mathrm{K}$. was found to be 47.313 E. U. from the band spectrum data.

The vapor pressure of solid and liquid carbon monoxide was measured and represented by the equations

```
Liquid \(68.09^{\circ}-83.13^{\circ} \mathrm{K}\).
    \(\log P(\mathrm{~cm})=.-\frac{477.3}{T}+11.23721-0.064129 T+2.5911 \times 10^{-4} T^{2}\)
Solid \(61.55^{\circ}-68.09^{\circ} \mathrm{K}\).
    \(\log P(\mathrm{~cm})=.-\frac{425.1}{T}+7.82259-0.0075960 T\)
```

Solid below $61.55^{\circ} \mathrm{K}$.

$$
\log P(\mathrm{~cm} .)=-\frac{418.2}{T}+4.127 \log T+1.47365-0.02623 T
$$

The free energy function for carbon monoxide was calculated from spectroscopic data and tabulated to $5000^{\circ} \mathrm{K}$. These data were combined with similar data on oxygen and with the available calorimetric data on graphite to give $\Delta F$ for the reaction $\mathrm{C}_{\text {graphite }}+1 / 2 \mathrm{O}_{2}=\mathrm{CO}$ to $3000^{\circ} \mathrm{K}$.

The free energy of formation of carbon monoxide from $\beta$ graphite and oxygen at $298.1^{\circ} \mathrm{K}$. was found to be $\Delta F_{298.1}^{\circ}=-33,000$.

Berkeley, California
[Contribution from the Department of Chemistry, Yale Univbrsity]

# THE ACTIVITY COEFFICIENT OF POTASSIUM CHLORIDE IN AQUEOUS SOLUTION FROM BOILING POINT DATA ${ }^{1}$ 

By Blair Saxton and Rodney P. Smith<br>Received January 18. 1932 Published July 6. 1932

Although the thermodynamic properties of aqueous salt solutions at ordinary temperatures have been extensively investigated, few data are available for the study of these properties at higher temperatures. Measurements of the boiling point elevation of solutions could be used for such a study if sufficient accuracy in the measurement could be attained. Considerable data ${ }^{2}$ have been published on the boiling point elevation of salt solutions. However, due to superheating, variation of pressure, errors of analysis or temperature measurement, the results are not of sufficient accuracy for a calculation of the activity coefficient. It has seemed worth while, therefore, to attempt to design an apparatus for a more accurate determination of the boiling point elevation.

## Apparatus

Of the different methods used for the elimination of superheating, the one suggested by Cottrell ${ }^{3}$ and used by Washburn and Read, ${ }^{4}$ Pearce and Hicks, ${ }^{5}$ and Bancroft and Davis ${ }^{6}$ seems to be the most satisfactory.

The first apparatus, Fig. 1, consisted of a Pyrex tube, A, 43 cm . long and 7.6 cm . inside diameter, over which was sealed a larger tube, $\mathrm{B}, 28 \mathrm{~cm}$. long, 11.5 cm . inside diameter. A condenser, $C$, and a small drain tube, $D$, were sealed into the outer tube, B. A tube, E, for the "cold" junction of the thermocouple was also sealed into the

[^5]
[^0]:    ${ }^{6}$ Moissan, Compt. rend., 102 (1886).
    ${ }^{7}$ Rayleigh, Proc. Roy. Soc. (London), A62, 204 (1897).
    ${ }^{8}$ Regnault, Ann. chim., 5, 52 (1842).
    ${ }^{9}$ Schlatter, Thesis, Geneva, 1923.
    ${ }^{10}$ Sternewarte, Landolt, Börnstein and Roth, "Physikalisch-cliemische Tabellen," Berlin, 1923.

[^1]:    ${ }^{1 s}$ Cailletet, Compt. rend., 85, (1877).
    ${ }^{14}$ Olszewsky, ibid., 100, 350 (1885).
    ${ }^{15}$ Von Wroblewsky. Wien. Ber., 90 (1885).
    ${ }^{16}$ Verschoyle, Trans. Roy. Soc. (London), A230, 189 (1931).

[^2]:    ${ }^{18}$ Clusius and Teske, Z. physik. Chem., [B] 6, 135 (1929).

[^3]:    ${ }^{19}$ Since this paper was written our attention has been called to some very accurate vapor pressure measurements by Verschoyle [Trans. Roy. Soc. (London), A230, 189 (1931)]. The equation given by Verschoyle does not represent his values in the liquid range, the discrepancy undoubtedly being due to some typographical error. However, we have compared his results from the melting point to $82.2^{\circ}$ with our Equation 2 . The agreement is extraordinarily good. The maximum deviation is $0.02^{\circ}$ and the average deviation $0.01^{\circ}$. However, Verschoyle's results, which extend to a pressure of 143 cm ., are not very satisfactorily represented by our Equation 2 above the boiling point. Below the melting point the agreement is not as good. Although Verschoyle's results extend to $54.21^{\circ} \mathrm{K}$., he apparently failed to notice the transition at $61.55^{\circ} \mathrm{K}$. and represents all of his results on the solid states by a single equation. Near the transition temperature Verschoyle's temperatures are about $0.06^{\circ}$ higher than the present work. Below $61.55^{\circ} \mathrm{K}$. Verschoyle's results do not agree well with our equation 3 obtained with the assistance of the calorimetric data. At $54.2^{\circ} \mathrm{K}$. lis results are $0.12^{\circ}$ lower than those given by Equation 3. However, our one observation on this crystal modification at $60.24^{\circ}$ is only $0.03^{\circ}$ lower than Verschoyle's results in this region.

    For comparison: Boiling point, $81.62^{\circ}$ (V.), $81.61^{\circ}$ (C. and G.).
    Triple point temperature, $68.06^{\circ}$ (V.), $68.09^{\circ}$ (C. and G.).
    Triple point pressure, 11.486 cm . (V.), 11.531 cm . (C. and G.).
    ${ }^{20}$ Von Wroblewsky, Compt. rend., 100, 979 (1885).

[^4]:    ${ }^{29}$ Roth and Naeser, Z. Elektrochem., 31, 461 (1925).
    ${ }^{30}$ Rossini, Bur. Standards J. Research, 6, 37 (1931).

[^5]:    ${ }^{1}$ Part of this paper is from a dissertation submitted by Rodney P. Smith to the Graduate School of Yale University, in partial ftifilment of the requirements for the degree of Doctor of Philosophy.
    ${ }^{2}$ For summary of results see "International Critical Tables," Vol. III. p. 324.
    ${ }^{3}$ Cottrell, This Journal, 41, 721 (1919).
    ${ }^{4}$ Washburn and Read, ibid., 41, 729 (1919).
    ${ }^{5}$ Pearce and Hicks, J. Phys. Chem., 30, 1678 (1926).
    ${ }^{6}$ Bancroft and Davis, ibid., 33, 591 (1929).

